

ION

has been devoted to  
of a wide range of so-  
be described as net-  
nternet [3, 4, 5, 6] or  
mmunities [8], food-  
networks [10, 11, 12,  
], in which nodes rep-  
nd links the physical  
y, many of these net-  
roperties and dynam-  
nted for by classical  
r, small-world prop-  
ributions [17] (where  
is defined as the num-  
ched) seem to emerge  
verning the topology.  
al properties imply a  
d a short average dis-  
nsiderable impact on  
aking place on top of  
e (SF) networks have  
om damage (absence  
[20] and prone to epi-  
shold) [21, 22, 23, 24].  
ological properties of  
l-world and scale-free  
on-trivial degree cor-  
Recently, an interest-  
duced by Klemm and

study how clustering the resilience to damage in networks [26, 27].

In this paper we re-examine the model. We find an alternative way to calculate the values of active sites. In addition, large scale numerical simulations show noticeable variability in the distribution of active sites. In particular, the distribution of active sites for the general case with a complex topology is also studied. The construction algorithm used in the numerical simulations we study is able to cover the whole range of parameters. We calculate the average number of active sites and connectivity corresponding to the variability with respect to the parameter  $m$ . Extensive numerical simulations allow us to get a clear picture presented here.

In the generated small-world properties we find a network of number of nodes forming topology is therefore star-shaped graphs. one dimensional lattice properties. In part cesses might be heard age distance among to a one-dimensional cuss the properties

with probability

$$p_d(K) = \frac{1}{1 + }$$

set of active nodes  $\mathcal{A}$ ,  
notes the in-degree of

del is quite sensitive  
3 are performed and,  
e the following cases.

before step 3.

after step 3.

e solved analytically  
on, after introducing  
ve node has in-degree  
order of steps 2 and  
n-degree distribution

(2)

using  $a = m$ . In this  
nversely proportional  
 $s(m+k^{\text{in}})^{-1}$  and the  
be  $P(k) = 2m^2k^{-3}$ .  
n mechanism the net-  
cient that approaches  
limit [25].

n claimed that finite  
ty distribution shows  
vior. We shall see in  
0 the model presents  
which yields a con-  
the  $m \rightarrow \infty$  limit. In  
ology is very sensible  
wing algorithms.

Each time the oldest  
increases by one and  
oldest node has in-deg-  
it is not deactivated  
degree 2, 3, ...,  $K - 1$  —  
creating a deactivation  
the probability that  
in  $K - 2$  steps and

$$\tilde{P}(K) =$$
  
$$=$$

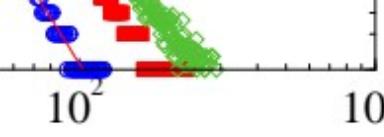
where  $\Gamma(x)$  is the s-  
the other hand, even if  
deactivated, the other  
Hence, in the  $K - 1$ -  
generation of a node with  
in-degree 1 are created  
with in-degree 1 cre-

$$\tilde{P}_1 = \sum_{K=1}^{\infty}$$

Therefore, the in-de-

$$P(k^{\text{in}}) = C$$

where  $C$  is a constan-  
condition  $\sum_{k^{\text{in}}} P(k^{\text{in}})$



A for  $a = m$ , network of  $m$ . The continuous  $\gamma = 2$  given by Eq. (10). exponent  $\gamma$  as a function of  $s$ .

previous expression us-  
ain that the in-degree  
behavior

$$= 2 + a. \quad (9)$$

If nodes is  $m$  then the out-degree) is  $m + k^{\text{in}}$  on shifted by  $m$ . For the degree distribution

$$k = 3 \quad (10)$$

$$\overline{3}) \quad k > 3$$

olution obtained from or  $a = m$ . For  $m = 2$  d agreement with the (0), with a power law . In the limiting case predicts the exponent er bound. Hence,

$$3 < \gamma \leq 4 \quad (11)$$

The probability that it has degree  $K$  is g

$$\tilde{P}(K) = \\ =$$

In the process of cre-  
nodes of in-degree  
number of nodes wi

$$\tilde{P}_0 = \sum_{K=1}^{\infty}$$

Thus, the analytic degree distribution :

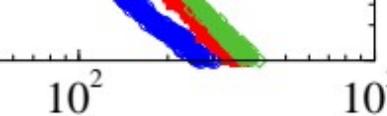
$$P(k^{\text{in}}) = C$$

with the normalizat

$$C = \tilde{P}_0$$

From here follows th  
tion (where  $k = m -$

$$P(k) = \begin{cases} \frac{\Gamma(2-a)}{\Gamma(a)} k^{-a}, & k \geq 3 \\ 0, & k < 3 \end{cases}$$



B for  $a = m$ , network of  $m$ . The inset shows on of  $m$  obtained from

ow a power law decay  
ntinuously increasing  
g that for  $m < 10$  the  
ngly differs from the

$n = 1$  the analytic so-  
readily seen from the

In fact, the solution  
n the thermodynamic  
weight of the nodes  
respect to the nodes  
singularity is rooted in  
an exponent  $\gamma = -2$ ,  
thermodynamic limit,  
model  $B$  has average  
ily implies that there  
the network size  $N$  in  
 $= 1$ , dependence that  
solution since we are  
work limit. We can  
ll form of the degree  
mposed by  $N$  nodes,  
 $k_c$ , such that there  
at  $k_c$ . Assuming that  
the same functional  
 $a = 1$

$= 1$

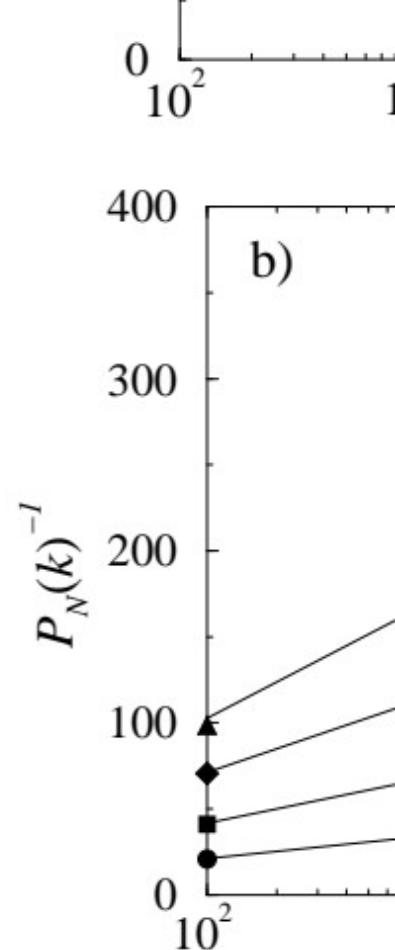


FIG. 3: Scaling of the model with  $a = m = 1$  size  $N$ , for (a)  $k = 1$  least-squares fits to the  $P_N(k)^{-1} \sim \ln N$  in (b)

proximation that re-

$$C_1 = 1 - \frac{2 \ln(3)}{\ln \left( \frac{1 + \sqrt{5}}{2} \right)}$$

For finite SF netwo-  
 $k^{-\gamma}$ , the maximum  
nodes as  $k_c \sim N^{1/\gamma}$   
 $k_c \sim N$ , and thus, fo-

$$1 - C_1$$

tion has a divergent

wn that the deactivation order in which steps 2 and 3 are carried out are distributions with a different behavior depending on the order. This is rather sensible to the range of variation. This previous works where this was done [27], prompting that in those works should

## EFFICIENT

tribution and compute the function of the node  $i$ . Then we can perform an average of  $a$  and  $m$  and for each node we compute the clustering coefficient of the network as undirected graph with total degree of the node

node  $i$  is defined by

(24)

between the neighbors of  $i$ . The maximum possible value of  $c_i(k)$  in our model new edges are added between  $i$  and the added node. Between active nodes and inactive nodes and in general, all the active nodes increase their degree. Moreover, all the active nodes increase their degree. At the same time we add a node to the network, the degree of  $i$  increases by one and the number of edges between  $i$  and its neighbors increases by one.

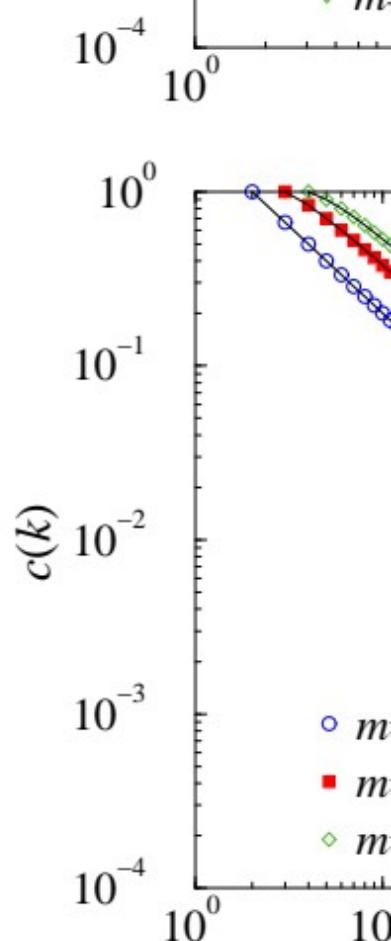


FIG. 4: Clustering coefficient  $c(k)$  for different values of  $m$ . Numerical simulations of a network of size  $N = 10^5$  and the analytical solution

$i$  was created. Besides, the initial degree of  $i$  is  $m$ , thus  $e_i(0) = m$ .  
1. Integrating Eq. (24) and substituting the result in Eq. (23), we find that  $t = k_i - m$ , we get

$$c(k) = \frac{m}{k} e^{-\frac{m}{k}} + \frac{2(m-1)}{k^2} e^{-\frac{2(m-1)}{k}}$$

where the last expression is obtained after some algebraic manipulations. The first term in the

s of the neighbors of the quantity  $k_{nn,i}$  does not depend on the degree of the node  $i$ . Such relations are present. The degree of the node  $i$  is increased. In particular, we consider the situation. In the first situation, the new node will connect more to nodes with high degrees; a property referred to as “hubs”. On the opposite side, it is called “leaving”; i.e. highly connected to nodes with low degrees.

When the node is added to the network,  $\langle k \rangle_A$ , where  $\langle k \rangle_A$  is the average degree of all neighbors. Then, if the node is a neighbor of the  $m - 1$  nodes, every time a new node is added, the average increases by  $(m - 1) + 1$ . This is due to the remaining  $m - 1$  nodes, the new one and the  $m$  because of the new connection.

$$+ m \langle k \rangle_A \quad (28)$$

Substituting this equation, taking  $D_i(0) = m \langle k \rangle_A$  and  $D_i(t) = D_i(0)$ , we get that

$$+ m \langle k \rangle_A \quad (29)$$

Now, when an active node remains fixed but the number of neighbors will still increase until reaching the infinite time limit,

$$+ m \langle k \rangle_A \quad (30)$$

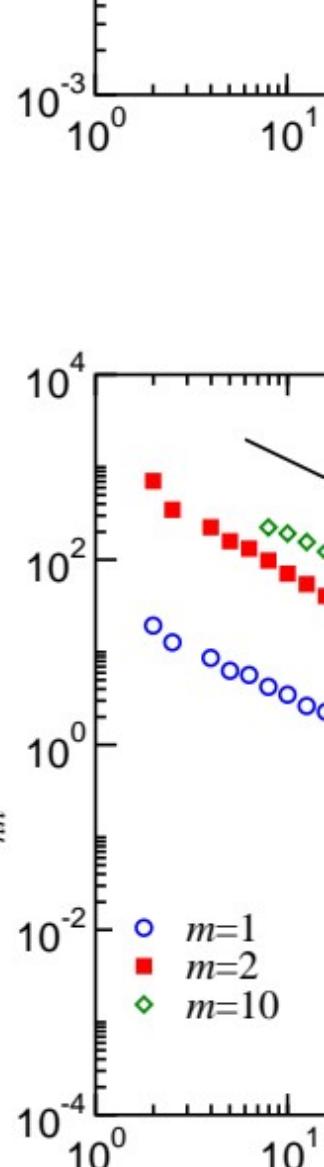
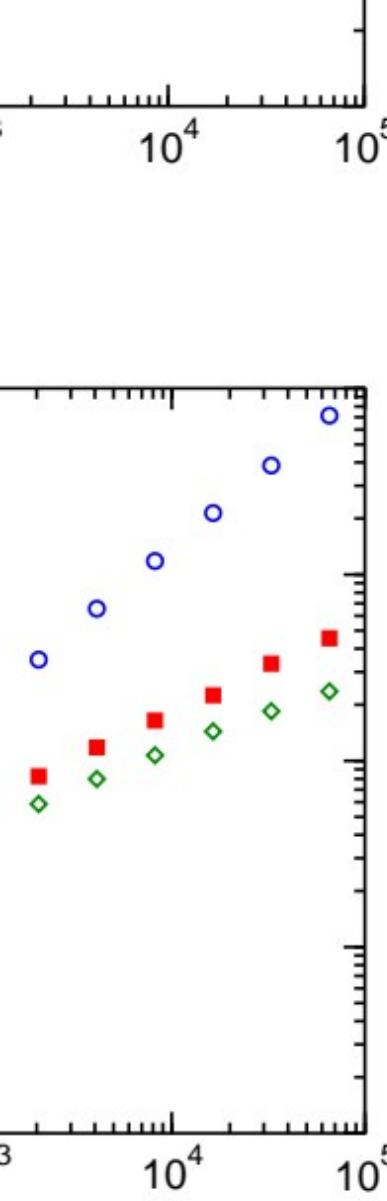


FIG. 5: Average near-degree  $\bar{k}_{nn}(k)$  for different  $m$  obtained from numerical simulations up to a network size of  $10^4$  nodes. The continuous line shows the dependency  $\bar{k}_{nn}(k) \sim k^{-\alpha}$ .

time average of  $\langle k \rangle_A$  is the connectivity of any dead node ( $m = 2$  for model I). The average degree of an active node is the sum of the degree of the remaining nodes having a degree  $k$  with  $m = 2$ , and depends only on  $k$  independently of the number of active nodes. Therefore, in this case, the



ree as a function of the of  $m$ . The points were of (a) model A and (b)  $5$ , averaging over 1000

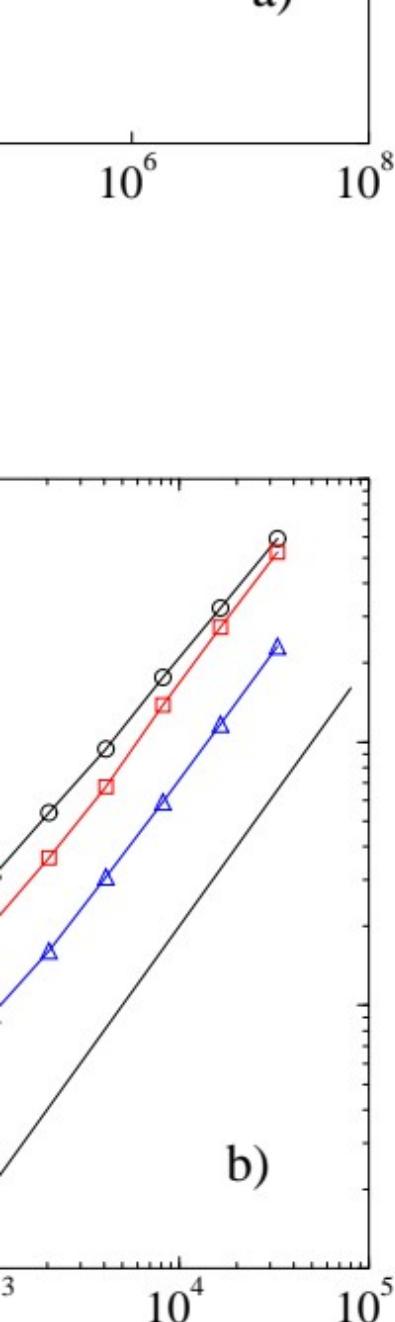
$1/k$  even for  $m \neq 2$ , the case of model B,  $m = 1$  and  $m = 10$  values of  $m$ . Thus, in the active node de- for intermediate val- we find that corre- e of “disassortative” s are preferably con- ss. It is also worth

increasing  $m$  the ex- limit  $\gamma = 3$  for  $m$  mically with  $N$ . O  $2 \leq \gamma < 3$ ,  $\langle \bar{k}_{nn} \rangle_N$  law. This implies th and that in the the form first the limit  $N$  connectivity curve i larger values. This nearest neighbor con quantity since the  $k$  after the  $N \rightarrow \infty$  li  $N$  is related to a ge diverging connectivit the detailed balance

## VI. DIAMETER

Another fundamental networks is identified length among nodes minimum path between imum number of in versed to go from n path length  $\langle d \rangle$  is th tance averaged over network. Similarly, the largest among nodes in the networ

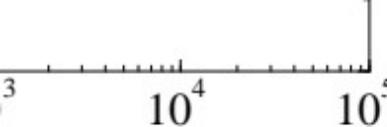
While regular netw tices) have a diamet inverse of the Euclid works show striking erage one can go fr the system by passi intermediate nodes



d the average shortest  
rent values of  $m$ . The  
ake of clarity, the curve  
factor 5.

implies that the topol-  
approaching those of a  
ords, the deactivation  
properties.  
sentation of the deac-  
in Fig. 8 the illustra-  
del B and  $a = m = 3$ .

FIG. 8: Illustration of the deactivation model B. The linear topology with  $a = m$  is evident.



random walker on the one dimensional lattice,  $N = 10^5$  nodes, as well as with 6301 nodes.

all discuss in the last properties might have properties of the net-

## CONCLUSIONS

provided a detailed analysis introduced in Ref. [25]. Being very sensible to the model and slight changes. The most striking feature is depending on the number of simultaneously active nodes  $m$  (i.e. when the deactivation rules' total degree. The approach to approaching the value properties of networks for  $1 \leq m \leq 10$ . Along in previous works, we testing degree correlations and marked disassortative connected nodes link degrees. The analytical expression is obtained and remarkably, the SF and

the percolation transition are extremely robust. A natural question is to what properties of SF networks are general results obtained. This reason, several effects of such correlations are important on these networks. In Ref. [26] it has been shown that the threshold in the case of B.

The presence of a deactivation model has been tracked by the gradient and the finite size effect on the connectivity of the lattice. We have shown here that the connectivity in the limit of large size. What appears to be the properties of spreading is a clear structure, with a chain-like structure. In a coarse grained limit, the system is dominated by the chain. In order to compare with a standard random walk, in Fig. 9 we plot the mean square displacement of a random walker,  $\langle R^2(t) \rangle$ . The brackets denote an average over a random walk on 250 nodes of a disassortative system, as would be the case for a lattice, we would expect the deactivation model to exhibit massive behavior, with a power law decay as  $\langle R^2(t) \rangle^{1/2} \sim t^{0.4}$ . However, on the deactivation model, the expected from its nature of spreading and the fact that the walk cannot therefore

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s the usual absence of  
pective, it would be

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